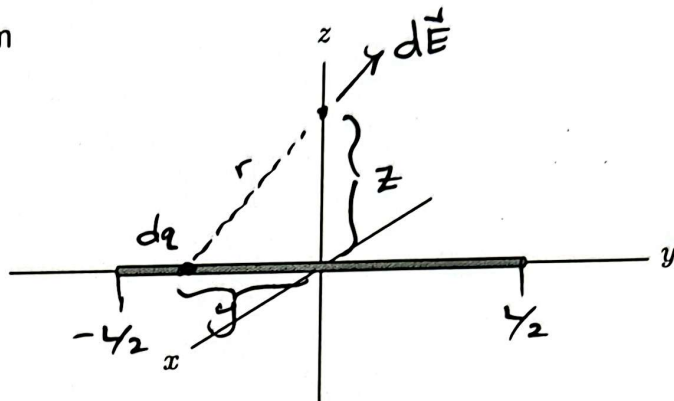


Find the electric field a distance z away from the center of a line charge (with charge density λ C/m).



$$\vec{E} = \int d\vec{E} = \int \frac{1}{4\pi\epsilon_0} \frac{dq}{r^2} \hat{r} = \frac{1}{4\pi\epsilon_0} \int \frac{dq}{r^2} \hat{r}$$

in terms of y :

$$dq = dy \cdot \lambda$$

$$r^2 = y^2 + z^2$$

$$\hat{r} = \hat{j} \sin\theta + \hat{k} \cos\theta$$

$$\sin\theta = \frac{-y}{\sqrt{y^2 + z^2}} \quad \cos\theta = \frac{z}{\sqrt{y^2 + z^2}}$$

$$\vec{E} = \frac{\lambda}{4\pi\epsilon_0} \int_{-L/2}^{L/2} \frac{dy}{y^2 + z^2} \cdot \left(-\hat{j} \frac{y}{\sqrt{y^2 + z^2}} + \hat{k} \frac{z}{\sqrt{y^2 + z^2}} \right)$$

drop this

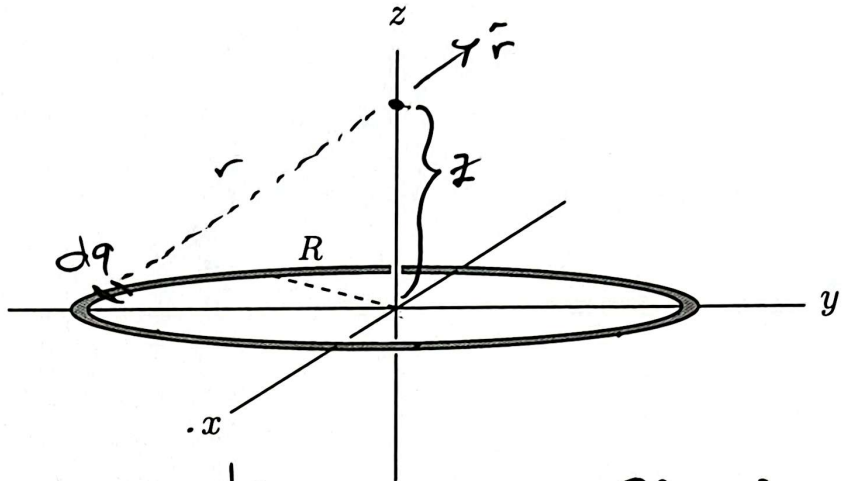
$$\vec{E} = \frac{\lambda}{4\pi\epsilon_0} \hat{k} z \int_{-L/2}^{L/2} \frac{dy}{(y^2 + z^2)^{3/2}} = \frac{\lambda}{4\pi\epsilon_0} z \hat{k} \left[\frac{y}{z^2 \sqrt{y^2 + z^2}} \right]_{-L/2}^{L/2}$$

$$= \frac{\lambda z \hat{k}}{4\pi\epsilon_0} \cdot \frac{L/2 \cdot 2}{z^2 \sqrt{(L/2)^2 + z^2}} = \frac{\lambda \hat{k} L}{4\pi\epsilon_0 z \sqrt{(L/2)^2 + z^2}}$$

Q = λL

$$= \frac{1}{4\pi\epsilon_0} \frac{Q}{z \sqrt{(L/2)^2 + z^2}} \hat{k}$$

Find the field along the z axis due to a ring of line charge (charge density λ C/m)

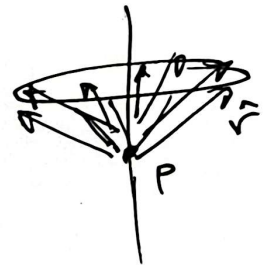
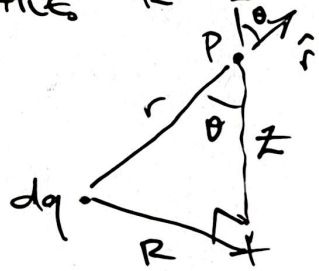


$$\vec{E} = \int d\vec{E} = \int \frac{1}{4\pi\epsilon_0} \frac{dq}{r^2} \hat{r}$$

$$= \frac{1}{4\pi\epsilon_0} \frac{1}{R^2+z^2} \int dq \cdot \hat{r}$$

$$r^2 = R^2 + z^2$$

(constant)



horizontal components cancel!

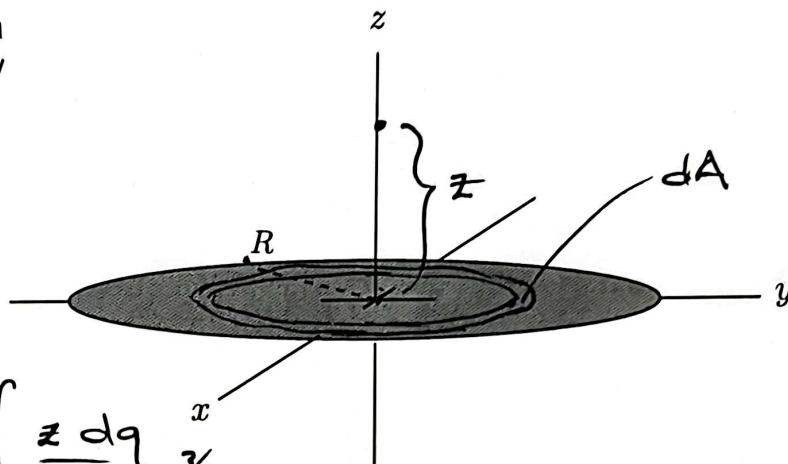
$$\hat{r} = (\cos\theta) \hat{k} + (\cancel{z}, \hat{j})$$

$$\cos\theta = \frac{z}{\sqrt{R^2+z^2}}$$

$$\vec{E} = \frac{1}{4\pi\epsilon_0} \frac{1}{(R^2+z^2)} \frac{z}{\sqrt{R^2+z^2}} \hat{k} \int dq$$

$$\vec{E} = \frac{1}{4\pi\epsilon_0} \frac{zQ}{(R^2+z^2)^{3/2}} \hat{k}$$

Find the field along the z axis due to a disk of surface charge (charge density σ C/m²)



$$\vec{E} = \int d\vec{E} = \frac{1}{4\pi\epsilon_0} \hat{k} \int \frac{z dq}{(r^2+z^2)^{3/2}}$$

field of a ring charge

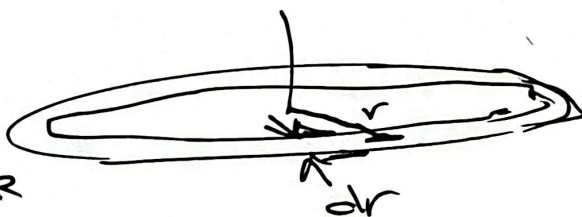
$$dq = \sigma dA$$

$$A = \pi r^2$$

$$\frac{dA}{dr} = 2\pi r$$

$$dq = \sigma \cdot 2\pi r \cdot dr$$

$$dA = 2\pi r \cdot dr$$



$$\vec{E} = \frac{1}{4\pi\epsilon_0} \hat{k} \cdot z \cdot \sigma \cdot 2\pi \int_0^R \frac{r dr}{(r^2+z^2)^{3/2}}$$

$$= \frac{1}{2\epsilon_0} \cdot z \sigma \hat{k} \cdot \left[-\frac{1}{\sqrt{r^2+z^2}} \right]_0^R$$

$$= \frac{z\sigma}{2\epsilon_0} \hat{k} \cdot \left(-\frac{1}{\sqrt{R^2+z^2}} + \frac{1}{z} \right)$$

$$\boxed{\vec{E} = \frac{\sigma}{2\epsilon_0} \left(1 - \frac{z}{\sqrt{R^2+z^2}} \right) \hat{k}}$$

How does this field behave when $z \gg R$?

$$\text{if } z \gg R \quad \sqrt{R^2 + z^2} \approx \sqrt{z^2} = z$$

$$\text{then } \vec{E} = \frac{\sigma}{2\epsilon_0} \left(1 - \frac{z}{z}\right) \hat{k} = 0$$

Sure, it goes to zero, but we want to know how it goes to zero.

To do this we can use the "binomial approximation", which says that

$$(1+x)^\alpha \approx 1 + \alpha x \quad \text{if } x \ll 1.$$

(This is really just the first two terms of the Taylor series.)

Consider the second term in the parentheses:

$$\frac{z}{\sqrt{R^2 + z^2}} \cdot \frac{1}{z} = \frac{1}{\sqrt{\left(\frac{R}{z}\right)^2 + 1}} = \left[1 + \left(\frac{R}{z}\right)^2\right]^{-1/2}$$

This has the form of the binomial above, with $x = \left(\frac{R}{z}\right)^2$. Clearly $\left(\frac{R}{z}\right)^2 \ll 1$. So:

$$\frac{z}{\sqrt{R^2 + z^2}} \approx \left(1 + \left(\frac{R}{z}\right)^2\right)^{-1/2} \approx 1 - \frac{1}{2} \left(\frac{R}{z}\right)^2$$

$$\vec{E} \approx \frac{\sigma}{2\epsilon_0} \left(1 - 1 + \frac{1}{2} \frac{R^2}{z^2}\right) \hat{k} = \frac{\sigma}{4\epsilon_0} \frac{R^2}{z^2} \hat{k}$$

$$\text{with } \sigma = \frac{Q}{\pi R^2}, \text{ this is } \vec{E} = \frac{Q}{4\pi\epsilon_0} \frac{1}{z^2} \hat{k} \quad \text{QED}$$